

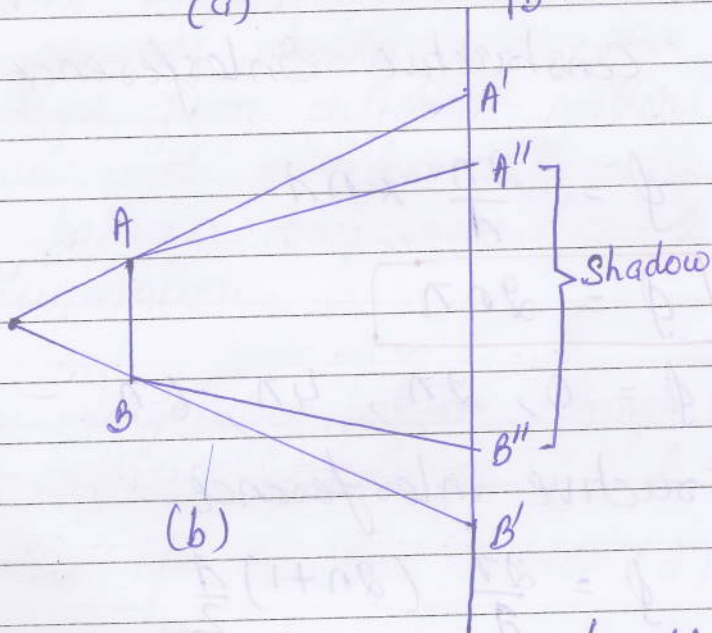
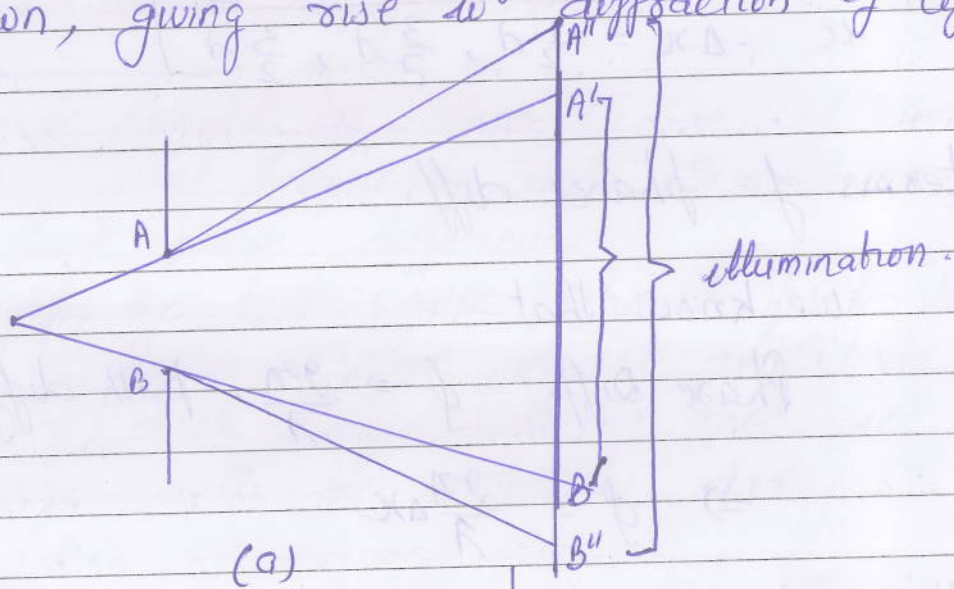
Diffraction of light \rightarrow

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Diffraction of light \rightarrow The phenomenon of bending of light around the corners of slit or obstacle is called diffraction.

When a slit or obstacle AB is placed in front of a light source then light bends around the corners of slit or obstacle as shown, giving rise to diffraction of light



Since light travels in the straight line, these position $A'B'$ should be illuminated (in fig a) or should have shadow (in fig b). But due to diffraction (bending of light), it is $A''B''$ instead

of radio & sound wave
 Diffraction is so common in our daily life but diffraction of light wave is not so common. Because for diffraction to take place, size of the slit should be of the order comparable to wavelength of wave. Since wavelength of sound waves is of the order of 6m and slits of size comparable to 6m are common w/o of size comparable to nano meters is so common. Thus diffraction of light is not common in our daily life.

Diffraction of light is of two types

1. **Fresnel Diffraction** \rightarrow When source of light and screen lies at a finite distance, then the diffraction that takes place is called Fresnel's diffraction. Since source is placed close to screen thus the incident wavefront is either spherical or cylindrical.

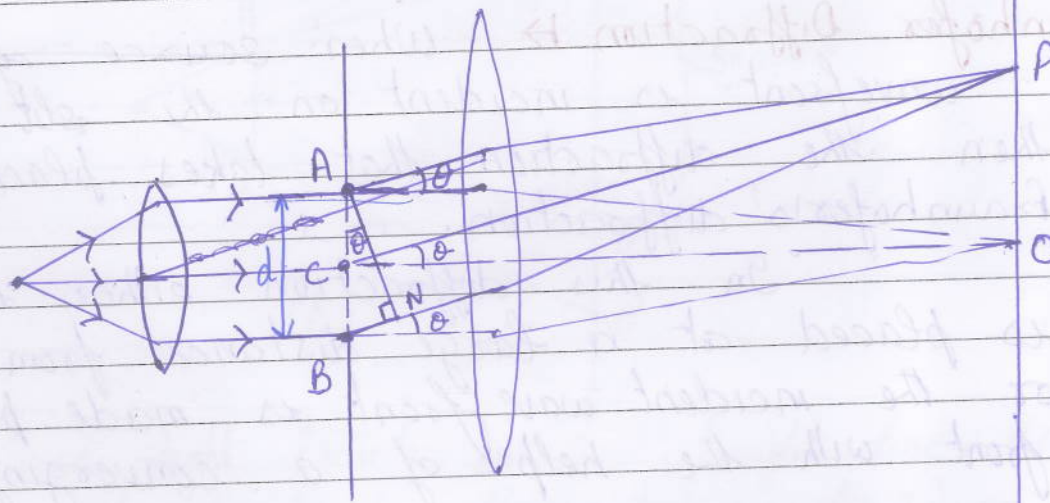
2. **Fraunhofer Diffraction** \rightarrow When source of a plane wavefront is incident on the slit or obstacle then the diffraction that takes place is called Fraunhofer's diffraction.

In this diffraction either the source is placed at a large distance from the slit or the incident wavefront is made plane wavefront with the help of a converging lens.

Diffraction at a single slit (Fraunhofer Diffraction)

Consider a plane wavefront from a monochromatic light is allowed to fall on a narrow slit AB ~~as shown~~ of width d as shown. According to Huygen's principle every point in b/w A and B acts as source of secondary disturbance. These secondary wavelets initially in same phase and reach at point O on the screen after travelling the same optical path and hence reinforce each other forming central maxima at point O .

Whereas the wavelets which get diffracted through different angles reach at various points on the screen and will have some path and hence phase difference. ~~Thus~~ Depending on the path difference, ~~various~~ secondary minimas or maximas are formed on screen.



Whereas for the secondary wavelets diffracted through angle θ , the path difference b/w the

wavelets from point A and B while reaching at point P will be BN (as optical path from A to P and NP to P is equal)

ie path difference, $BN = AB \sin \theta$ ($\because \frac{BN}{AB} = \sin \theta$)

$\Rightarrow \boxed{BN = d \sin \theta}$

Thus for different diffraction angles θ , path difference b/w the two wavelets from point A and B will be different. Thus various secondary maxima and minima will be formed on the various points of screen.

Suppose for diffraction angle θ_1 , the path diff. b/w the two wavelets from point A and B of the slit is d ,

ie $d \sin \theta_1 = d$,

then point P on the screen will be the position of 1st secondary minima. Because, the slit can be supposed to be divided into two equal parts AC and CB, so that for every point in the upper half AC there is a point in lower half CB having path diff of $d/2$ and phase difference of π . Thus these wavelets interfere destructively at point P and give rise to 1st secondary minima.

Similarly for diffraction angle θ_2 and path difference of $2d$

ie $d \sin \theta_2 = 2d$,

Again because, now the slit can be supposed to be divided into four equal parts. having path difference of $\lambda/2$ with the neighbouring part. Again wavelets from four parts interfere destructively producing 2nd secondary minima.

Wly for $d \sin \theta_n = n \lambda$, we get n^{th} secondary minima.

Suppose for diffraction angle θ_1' , the path difference b/w the two wavelets from point A and B is $\frac{3}{2} \lambda$

$$d \sin \theta_1' = \frac{3}{2} \lambda$$

Here point P' will be the position of 1st secondary maxima. Because, now the slit can be supposed to be divided into 3 equal parts having path difference of $\lambda/2$ with the neighbouring part. Thus the wavelets from the two parts having path diff. of $\lambda/2$ cancel out each other w/a wavelet from 3rd part reinforce each other producing 1st secondary ~~minima~~ ^{maxima} of comparably lower intensity.

Similarly for $d \sin \theta_2' = \frac{5}{2} \lambda$

we will get second ~~or~~ secondary minima. Because now the slit can be supposed to be divided into 5 equal parts. and light from four ^{CLASSMATE} parts cancel out each other and we get PAGE

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light of comparative lower intensity from the 5th part of the slit.

illy for $d \sin \theta_n' = (2n+1) \frac{\lambda}{2}$ we will get nth secondary maxima.

Intensity of secondary maxima goes on decreasing because for 1st secondary maxima we get light from $\frac{1}{3}$ rd of slit and for 2nd secondary maxima from $\frac{1}{5}$ th and for 3rd secondary maxima from $\frac{1}{7}$ th part of the slit and so on.

Intensity of secondary maxima relative to the intensity of central maxima are in the ratio of $1 : \frac{1}{9} : \frac{1}{25} : \frac{1}{49} \dots$

Width of secondary maxima, minima and central maxima \rightarrow

Suppose P is position of nth secondary minima.

Then $d \sin \theta_n = n \lambda$

$$\Rightarrow \sin \theta_n = \frac{n \lambda}{d} \quad \text{--- (1)}$$

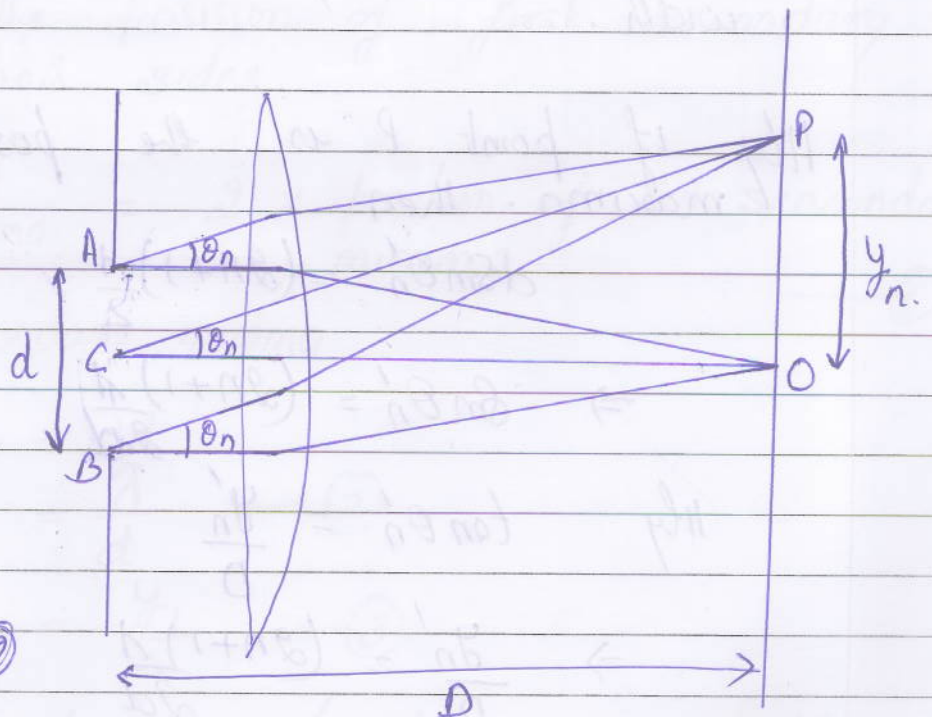
Also in ΔCOP

$$\tan \theta_n = \frac{y_n}{D} \quad \text{--- (2)}$$

For θ to be very small

classmate

$$\sin \theta_n \cong \tan \theta_n$$



$$\Rightarrow \frac{y_n}{D} = \frac{n\lambda}{d}$$

[using ① and ②]

$$\Rightarrow y_n = \frac{nD\lambda}{d}$$

Since secondary maxima lie b/w two successive minimas. Thus width the secondary maxima

$$\begin{aligned} \Rightarrow \beta &= y_n - y_{n-1} \\ &= \frac{nD\lambda}{d} - \frac{(n-1)D\lambda}{d} \\ &= \frac{D\lambda}{d} (n - n + 1) \end{aligned}$$

$$\boxed{\beta = \frac{D\lambda}{d}}$$

Since β is independent of n , therefore all the secondary maxima are of same width.

lly if point P is the position of n^{th} secondary maxima. then

$$d \sin \theta'_n = (2n+1) \frac{\lambda}{2}$$

$$\Rightarrow \sin \theta'_n = \frac{(2n+1)\lambda}{2d}$$

$$\text{lly } \tan \theta'_n = \frac{y'_n}{D}$$

$$\Rightarrow \frac{y'_n}{D} = \frac{(2n+1)\lambda}{2d}$$

{ $\because \sin \theta \approx \tan \theta$ }

$$\Rightarrow y_n' = (2n+1) \frac{\Delta D}{2d}$$

Thus width of secondary minima is

$$\begin{aligned} \beta' &= y_n' - y_{n-1}' \\ &= (2n+1) \frac{\Delta D}{2d} - [2(n-1)+1] \frac{\Delta D}{2d} \\ &= \frac{\Delta D}{2d} [2n+1 - 2n+2-1] \end{aligned}$$

$$\boxed{\beta' = \frac{\Delta D}{d}}$$

$$\text{i.e. } \beta = \beta'$$

~~Also~~ Thus all secondary maxima and minima are of same width.

Width of central maxima \rightarrow The central maxima extends upto to the position of first secondary minima of both sides.

Thus width of central maxima = 2 x position of 1st secondary minima. — (1)

for 1st secondary minima

$$d \sin \theta_1 = \lambda$$

$$\sin \theta_1 = \frac{\lambda}{d} \quad \text{--- (2)}$$

$$\text{also } \tan \theta_1 = \frac{y_1}{D} \quad \text{--- (3)}$$

but θ for θ to be very small

$$\sin \theta_1 \cong \tan \theta_1$$

$$\Rightarrow \frac{\lambda}{d} = \frac{y_1}{D}$$

$$\Rightarrow y_1 = \frac{\lambda D}{d}$$

$$\therefore \text{Width of Central maxima} = 2y_1 \quad \text{using (1)}$$

$$\Rightarrow \boxed{\text{Width of central maxima} = \frac{2\lambda D}{d}}$$

$$\text{or width of central maxima} = 2\beta \text{ or } 2\beta'$$

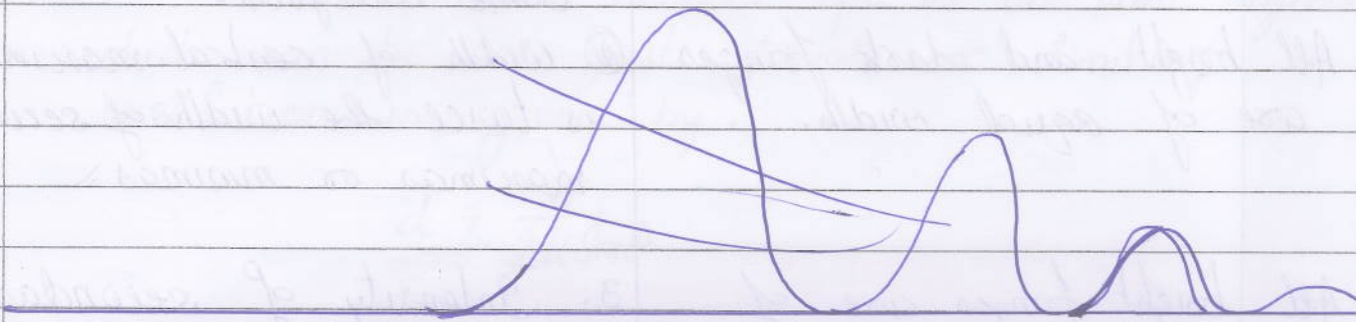
Thus width of central maxima is 2 times the width of 1st secondary maxima or minima w/a intensity of central maxima is 21 times the intensity 1st secondary maxima and 61 times the intensity of 2nd secondary maxima.

Note: \rightarrow

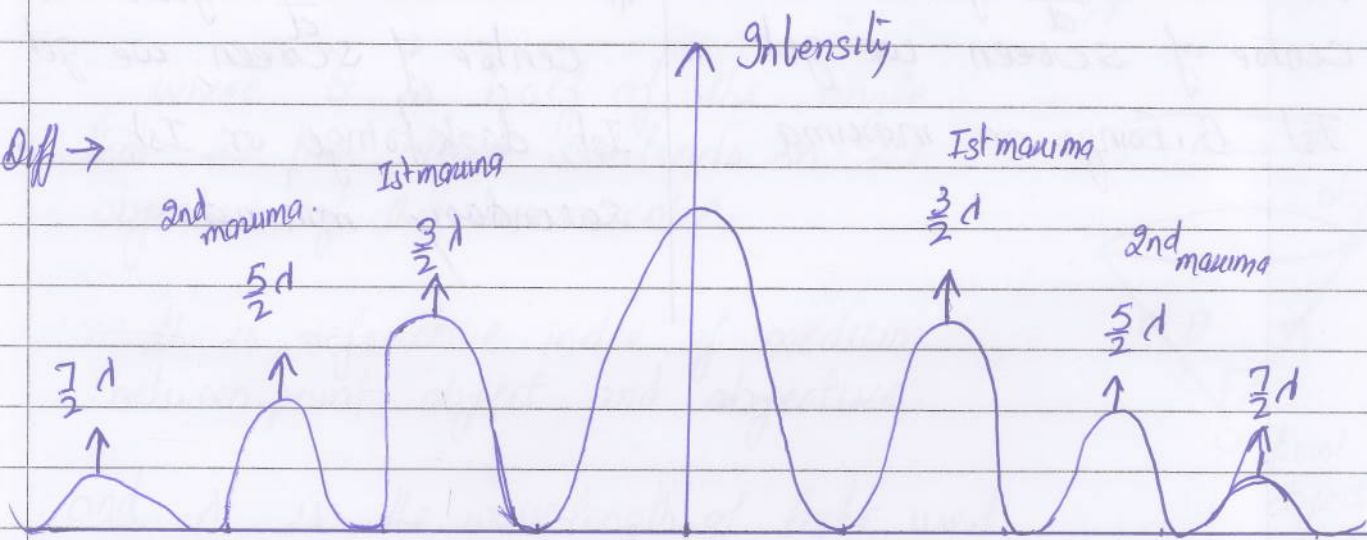
$$\text{Angular width of Central maxima} = 2\theta_1 = \frac{2\lambda}{d}$$

$$\text{Angular width of Secondary maxima or minima} = \frac{\lambda}{d}$$

Intensity Distribution curve \rightarrow The graph b/w intensities of maxima and/minimas and diffraction angle is called intensity distribution curve.



Path Diff \rightarrow



$\theta \rightarrow$	$-\frac{3\lambda}{d}$	$-\frac{2\lambda}{d}$	$-\frac{\lambda}{d}$	0	$\frac{\lambda}{d}$	$\frac{2\lambda}{d}$	$\frac{3\lambda}{d}$
$y \rightarrow$	$\frac{3D\lambda}{d}$	$\frac{2D\lambda}{d}$	$\frac{D\lambda}{d}$	0	$\frac{D\lambda}{d}$	$\frac{2D\lambda}{d}$	$\frac{3D\lambda}{d}$
Path diff.	3 λ	2 λ	1 λ		1 λ	2 λ	3 λ
	3rd minima	2nd minima	1st minima		1st minima	2nd minima	3rd minima

Difference b/w interference and diffractionInterference

- ① Interference is due to superposition of secondary wavelets from two diff. coherent sources
2. All bright and dark fringes are of equal width.
3. All bright fringes are of same intensity.
4. At distance $\frac{D\lambda}{d}$ from the centre of screen we get 1st B. fringe or maxima

Diffraction

Diffraction is due to scapes position of secondary wave from the different parts of same wavefront.

- ① width of central maxima is twice the width of secondary maxima or minima.
3. Intensity of secondary maxima goes on decreasing with order.
- ④ At distance $\frac{D\lambda}{d}$ from the centre of screen, we get 1st dark fringe or 1st secondary minima.

Resolving power of microscope and telescope \rightarrow

Resolving power of the microscope \rightarrow It is d/a the reciprocal of the smallest distance b/w two point objects at which they can be seen just separate through a microscope.

The smallest distance b/w the two point objects which can be seen just separate or limit of resolution is given by,

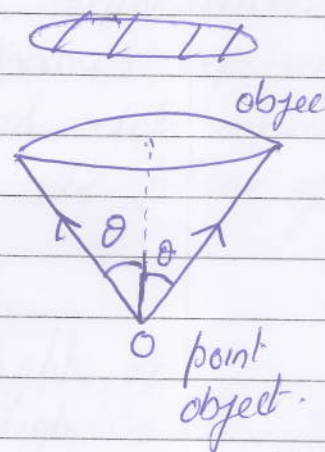
$$d = \frac{\lambda}{2\mu\sin\theta}$$

$$\begin{aligned} \text{Thus resolving power of microscope} &= \frac{1}{d} \\ &= \frac{2\mu\sin\theta}{\lambda} \end{aligned}$$

where θ is half of the angle that the point object subtends on the objective of the microscope.

μ is refractive index of medium between point object and objective.

and λ is the wave-length of light used.



Thus to increase resolving power of the microscope, oil of high value of refractive index is placed between objective of the microscope and point object.

Also resolving power of microscope will be high when light of smaller wave length (ie blue light) is used to illuminate the object.

Resolving power of the telescope \rightarrow It is d/a the smallest angular separation b/w the two distant objects which can be seen just separate through a telescope

The smallest angular separation or limit of resolution is given as

$$d\theta = \frac{1.22\lambda}{D}$$

$$\begin{aligned} \text{Thus Resolving power of telescope} &= \frac{1}{d\theta} \\ &= \frac{D}{1.22\lambda} \end{aligned}$$

where D is ~~apex~~ aperture or diameter of the objective of telescope

